

On an unsymmetric eigenvalue problem governing free vibrations of fluid-solid structures

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This is joint work with Heinrich Voss

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- 1 Problem definition and properties
- 2 Numerical methods
- 3 Numerical Results

Outline

1 Problem definition and properties

2 Numerical methods

3 Numerical Results

Problem definition

Vibrations of fluid-solid structures can be modelled in terms of solid displacement and fluid pressure and one obtains the classical form of an eigenproblem.

$$\operatorname{Div} \sigma(u) + \lambda \rho_s u = 0 \text{ in } \Omega_s,$$

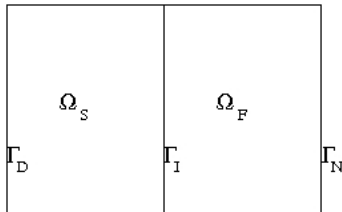
$$\Delta p + \frac{\lambda}{c^2} p = 0 \text{ in } \Omega_f,$$

$$\sigma(u) n - p n = 0 \text{ on } \Gamma_I,$$

$$\nabla p n - \lambda \rho_f u n = 0 \text{ on } \Gamma_I,$$

$$u = 0 \text{ on } \Gamma_D,$$

$$\nabla p n = 0 \text{ on } \Gamma_N,$$



where

- u : solid displacement
- p : fluid pressure
- λ : eigenparameter
- $\sigma(u)$: linearized stress tensor
- ρ_s, ρ_f : densities of solid and fluid

Problem definition, cntd.

This eigenvalue problem can be given an unsymmetric variational formulation which can be discretized by the Finite-Element method and one obtains the unsymmetric matrix eigenproblem

$$KX := \begin{pmatrix} K_s & C \\ 0 & K_f \end{pmatrix} \begin{pmatrix} x_s \\ x_f \end{pmatrix} = \lambda \begin{pmatrix} M_s & 0 \\ -C^T & M_f \end{pmatrix} \begin{pmatrix} x_s \\ x_f \end{pmatrix} =: \lambda M X, \quad (1)$$

where

- $K_s, M_s \in \mathbb{R}^{s \times s}$ are symmetric positive definite stiffness and mass matrices of the solid,
- $K_f, M_f \in \mathbb{R}^{f \times f}$ are symmetric stiffness and mass matrices of the fluid,
where K_f is semi positive-definite and M_f positive definite,
- $C \in \mathbb{R}^{s \times f}$ is due to the coupling effects between fluid and solid,
- $x_s \in \mathbb{R}^s$ is the solid displacement vector, and
- $x_f \in \mathbb{R}^f$ the fluid pressure vector.

This talk considers the properties of eigenproblem (1) and discusses ways how to use the symmetry of K_s, K_f, M_s , and M_f to adapt symmetric eigensolvers to the given problem.

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Properties

Some properties can easily be derived:

Lemma

- (1) can be symmetrized by $T := \begin{pmatrix} M_s^{-1} K_s & M_s^{-1} C \\ 0 & I \end{pmatrix}$, i.e.

$$T^T K x = \lambda T^T M x$$

is a symmetric eigenvalue problem.

- (1) has only real non-negative eigenvalues.
- If $x := \begin{pmatrix} x_s \\ x_f \end{pmatrix}$ is a right eigenvector of (1) corresponding to the eigenvalue λ , then $\hat{x} := \begin{pmatrix} \lambda x_s \\ x_f \end{pmatrix}$ is a left eigenvector.
- Right eigenvectors can be chosen orthonormal with respect to $\tilde{M} := \begin{pmatrix} K_s & 0 \\ 0 & M_f \end{pmatrix}$,
left eigenvectors can be chosen orthogonal with respect to $\bar{M} := \begin{pmatrix} M_s & 0 \\ 0 & K_f \end{pmatrix}$.
- Right eigenvectors x and left eigenvectors \hat{x} corresponding to distinct eigenvalues satisfy $\hat{x}^T K x = \hat{x}^T M x = 0$.

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Properties, cntd.

Lemma

Let $\lambda_j(A, B)$ denote the j smallest eigenvalue of the eigenproblem $Ax = \lambda Bx$ regarding the multiplicity of eigenvalues. Then it holds that

$$\begin{aligned}\lambda_j(K, M) &\leq \lambda_j(K_s, M_s), \quad j = 1, \dots, s \\ \lambda_{s+f+1-j}(K, M) &\geq \lambda_{s+1-j}(K_s, M_s), \quad j = 1, \dots, s \\ \lambda_j(K, M) &\leq \lambda_j(K_f, M_f), \quad j = 1, \dots, f \\ \lambda_{s+f+1-j}(K, M) &\geq \lambda_{f+1-j}(K_f, M_f), \quad j = 1, \dots, f.\end{aligned}$$

Proof:

Let $E_s := \text{span}\{e_1, \dots, e_s\}$ where $e_j \in \mathbb{R}^{s+f}$ denotes the j th unit vector containing a 1 in its j th component and zeros elsewhere. Then it holds that

$$\begin{aligned}\lambda_j(K, M) &= \min_{\dim V=j} \max_{x \in V, x \neq 0} \frac{x^T T^T K x}{x^T T^T M x} \leq \min_{\dim V=j, V \subset E_s} \max_{x \in V, x \neq 0} \frac{x^T T^T K x}{x^T T^T M x} \\ &= \min_{\dim W=j, W \subset \mathbb{R}^s} \max_{y \in W, y \neq 0} \frac{y^T K_s M_s^{-1} K_s y}{y^T K_s y} = \lambda_j(K_s, M_s).\end{aligned}$$

The second inequality is obtained analogously from the maxmin characterization, and the third and fourth inequalities follow in a similar way.

An inverse-free Rayleigh functional

A Rayleigh quotient for fluid-solid eigenproblems is given immediately by its symmetrized version. As it involves inverse matrices it is numerically less valuable and we are interested in an inverse-free analogon.

For a given right eigenvector $\begin{pmatrix} x_s \\ x_f \end{pmatrix}$ corresponding to the eigenvalue λ it holds

$$\lambda = \frac{\begin{pmatrix} \lambda x_s \\ x_f \end{pmatrix}^T \begin{pmatrix} K_s & C \\ 0 & K_f \end{pmatrix} \begin{pmatrix} x_s \\ x_f \end{pmatrix}}{\begin{pmatrix} \lambda x_s \\ x_f \end{pmatrix}^T \begin{pmatrix} M_s & 0 \\ -C^T & M_f \end{pmatrix} \begin{pmatrix} x_s \\ x_f \end{pmatrix}}$$

This suggests to define a Rayleigh functional ρ for some general $s + f$ -dimensional vector by the requirement

$$\rho(x_s, x_f) = \frac{\rho(x_s, x_f) x_s^T K_s x_s + \rho(x_s, x_f) x_s^T C x_f + x_f^T K_f x_f}{\rho(x_s, x_f) x_s^T M_s x_s - x_f^T C^T x_s + x_f^T M_f x_f}$$

which leads to

$$\rho(x_s, x_f)^2 x_s^T M_s x_s + \rho(x_s, x_f) (x_f^T M_f x_f - x_s^T K_s x_s - 2x_s^T C x_f) - x_f^T K_f x_f = 0.$$

We therefore choose the unique positive root of this equation as Rayleigh functional.

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This suggests to define a Rayleigh functional p for some general $s + f$ -dimensional vector by the requirement

$$p(x_s, x_f) = \frac{p(x_s, x_f) x_s^T K_s x_s + p(x_s, x_f) x_s^T C x_f + x_f^T K_f x_f}{p(x_s, x_f) x_s^T M_s x_s - x_f^T C^T x_s + x_f^T M_f x_f}$$

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Rayleigh functional and properties

Definition

$$p(x_s, x_f) := \begin{cases} q(x_s, x_f) + \sqrt{q(x_s, x_f)^2 + \frac{x_f^T K_f x_f}{x_s^T M_s x_s}} & \text{if } x_s \neq 0 \\ \frac{x_f^T K_f x_f}{x_f^T M_f x_f} & \text{if } x_s = 0 \end{cases} \quad (2)$$

where

$$q(x_s, x_f) := \frac{x_s^T K_s x_s - x_f^T M_f x_f + 2x_s^T C x_f}{2x_s^T M_s x_s}$$

is called Rayleigh functional of the fluid-solid vibration eigenvalue problem (1).

Well-known properties for symmetric eigenproblems can be generalized to the fluid-solid eigenproblem using the Rayleigh functional (2):

Lemma

Any eigenvector x of (1) is a stationary point of p , i.e. $\nabla p(x) = 0$.

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Rayleigh functional and properties, cntd.

As for symmetric eigenproblems one can derive variational formulations.

Proposition

The following variational characterizations hold

i) (Minmax characterization)

$$\lambda_k = \min_{\dim S_k = k} \max_{0 \neq x \in S_k} \rho(x) = \max_{\dim S_k = n+1-k} \min_{0 \neq x \in S_k} \rho(x).$$

ii) (Rayleigh's principle)

$$\begin{aligned} \lambda_k &= \min\{\rho(x) : x^T \tilde{M} x_j = 0, j = 1, \dots, k-1\} \\ &= \max\{\rho(x) : x^T \tilde{M} x_j = 0, j = k+1, \dots, s+1\} \end{aligned}$$

These variational characterizations can be proved by the following lemma.

Lemma

Assume that $x = \sum_{i \in I} \alpha_i x_i$ is a linear combination of some eigenvectors indexed by I . Then the Rayleigh functional is bounded by

$$\min_{i \in I} \lambda_i \leq \rho(x) \leq \max_{i \in I} \lambda_i.$$

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Rayleigh functional iteration

This Rayleigh functional can be used to define a Rayleigh functional iteration:

Require: initial vector $x^{(0)}$, $k = 0$

1: **repeat**

2: $\rho^{(k)} = \rho(x^{(k)})$

3: $x^{(k+1)} = (K - \rho^{(k)}M)^{-1} Mx^{(k)}$

4: Normalize $x^{(k+1)}$ by \tilde{M}

5: $k \leftarrow k + 1$

6: **until** convergence

As done for the Rayleigh quotient iteration for symmetric eigenproblems, local cubical convergence can be shown.

Proposition

The iterates $\rho^{(k)}$ and $x^{(k)}$ converge locally cubically towards to an eigenvalue λ and a corresponding eigenvector x .

Rayleigh functional iteration converges fast, but is highly sensitive with respect to the given initial vector. To overcome this drawback, one can consider projection methods where the projection space V is extended by the direction of the Rayleigh functional iteration. To reduce the computational effort we wish to solve

$$(K - \rho^{(k)}M)x^{(k+1)} = Mx^{(k)} \quad (3)$$

only approximately by some iterative method.

Jacobi-Davidson method

However, close to the desired eigenvector, the expanded subspace is very sensitive to inexact solves of (3) and hence we will replace the expansion direction by $t := x^{(k)} + \alpha x^{(k+1)}$ with α determined such that $x^{(k)T} t = 0$ which leads to the equivalent so called correction equation

$$\left(I - \frac{Mxx^T}{x^T Mx}\right)(K - \theta M)\left(I - \frac{xx^T}{x^T x}\right)t = -(K - \theta M)x, \quad t^T x = 0$$

for a given eigenpair approximation (θ, x) . Iterative projection methods of this type were introduced by Sleipen and van der Vorst and are known as Jacobi-Davidson method. For fluid-solid eigenproblems we still have to ensure real eigenvalues of the projected eigenproblem and therefore we use structure preserving projection matrices

$$V = \begin{pmatrix} V_s & 0 \\ 0 & V_f \end{pmatrix}$$

and obtain

$$V^T K V = \begin{pmatrix} V_s^T K_s V_s & V_s^T C V_f \\ 0 & V_f^T K_f V_f \end{pmatrix} \quad \text{and} \quad V^T M V = \begin{pmatrix} V_s^T M_s V_s & 0 \\ -V_f^T C^T V_s & V_f^T M_f V_f \end{pmatrix}.$$

Proposition

Assume that V has maximal rank l . Then it holds

$$\lambda_k(K, M) \leq \lambda_k(V^T K V, V^T M V) \quad \text{for } k = 1, \dots, l.$$

Jacobi-Davidson method, cntd.

Require: Initial basis $V = \begin{pmatrix} V_s & 0 \\ 0 & V_f \end{pmatrix}$, $V_s^T V_s = V_f^T V_f = I$, $m = 1$

- 1: determine preconditioner $L \approx (K - \sigma M)^{-1}$, σ close to the first desired eigenvalue
- 2: **while** $m \leq$ number of wanted eigenvalues **do**
- 3: compute the m smallest eigenvalue θ_m and corresponding eigenvector $y = \begin{pmatrix} y_s^T & y_f^T \end{pmatrix}^T$ of the projected problem

$$\begin{pmatrix} V_s^T K_s V_s & V_s^T C V_f \\ 0 & V_f^T K_f V_f \end{pmatrix} \begin{pmatrix} y_s \\ y_f \end{pmatrix} = \theta \begin{pmatrix} V_s^T M_s V_s & 0 \\ -V_f^T C V_s & V_f^T M_f V_f \end{pmatrix} \begin{pmatrix} y_s \\ y_f \end{pmatrix}.$$

- 4: determine Ritz vector $x = \begin{pmatrix} V_s y_s \\ V_f y_f \end{pmatrix}$ and residual $r = (K - \theta_m M)x$

5: **if** $\|r\|/\|x\| < \epsilon$ **then**

- 6: accept approximate m th eigenpair (θ_m, x) ; increase $m \leftarrow m + 1$; reduce search space
- 7: determine new preconditioner $L \approx (K - \theta_m M)^{-1}$ if necessary
- 8: choose approximation (θ_m, x) to next eigenpair
- 9: compute residual $r = (K - \theta_m M)x$

10: **end if**

- 11: compute approximate solution $t = \begin{pmatrix} t_s^T & t_f^T \end{pmatrix}^T$ of the correction equation

$$\left(I - \frac{Mxx^T}{x^T Mx}\right)(K - \theta M)\left(I - \frac{xx^T}{x^T x}\right)t = r, x^T t = 0$$

- 12: orthogonalize $v_s = t_s - V_s V_s^T t_s$, $v_f = t_f - V_f V_f^T t_f$
- 13: expand search space $V_s \leftarrow [V_s, v_s/\|v_s\|]$, $V_f \leftarrow [V_f, v_f/\|v_f\|]$ and update proj. problem
- 14: **end while**

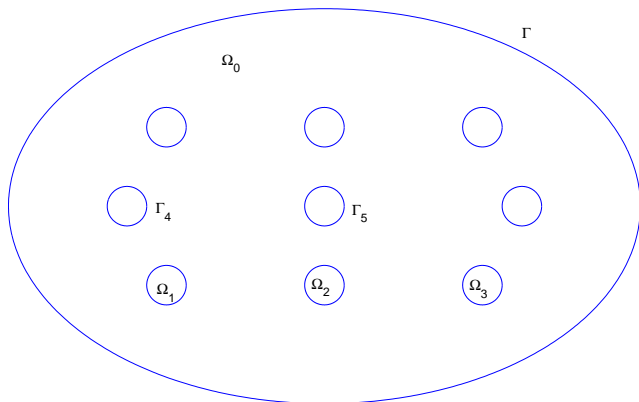
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Numerical Results

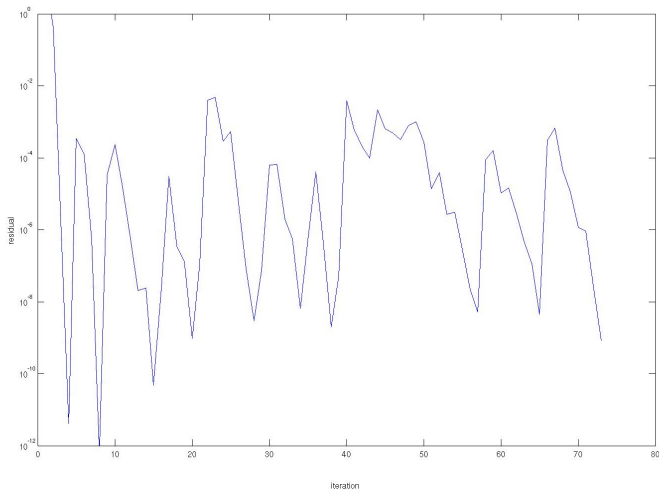
We consider the 2D-example of 9 tubes immersed in a fluid. The model has been discretized by the FEM with 9162 degrees of freedom. The task is to find the 10 smallest eigenvalues by the Jacobi-Davidson method.

Ω_0 : section of cavity, Ω_j : section of tube j



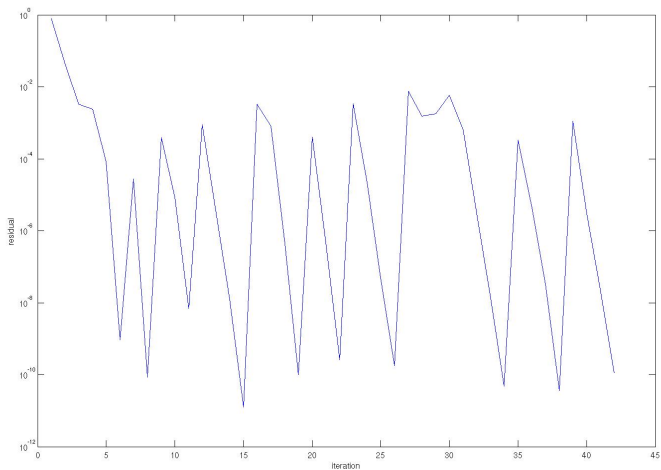
Numerical Results: JD method for general unsymmetric eigenproblems

Convergence history for Jacobi-Davidson method for general nonsymmetric eigenproblems



required iterations for 10 smallest eigenvalues: 73

Convergence history for fluid-solid interaction Jacobi-Davidson method



required iterations for 10 smallest eigenvalues: 42